

Hartree-Fock phase diagram of the Homogeneous Electron Gas in 2 or 3 Dimensions

L. Baguet, directed by B. Bernu

Collaborating with :
F. Delyon (CPHT,X)

M. Holzmann (LPTMC, Paris and LPMMC, Grenoble)

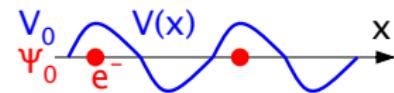
LPTMC, Université Paris VI

28 Novembre 2013

Why Hartree-Fock is still interesting ?

For many body systems : mean field must be done first. But Hartree-Fock ground state of HEG is still unknown !

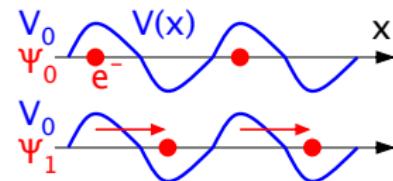
Iterative Method (Lanczos) for homogeneous system : bi-stable, not converging in general.



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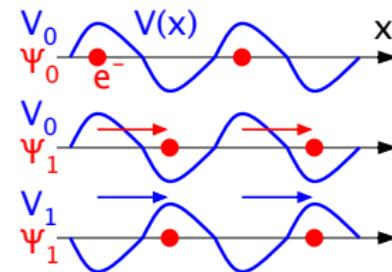
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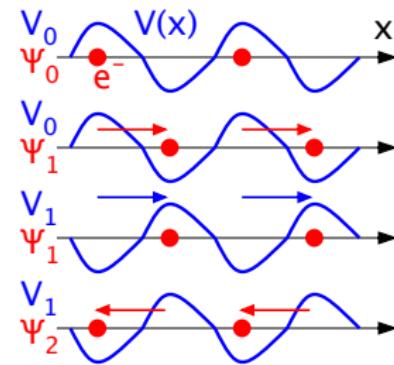
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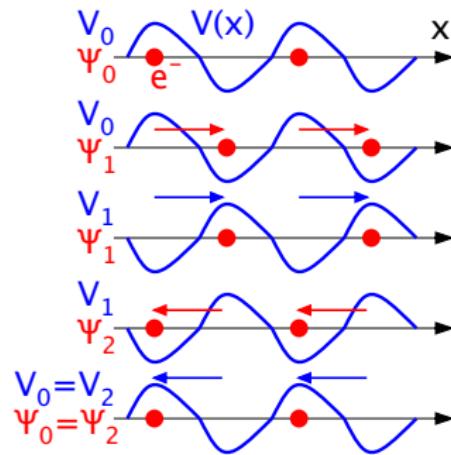


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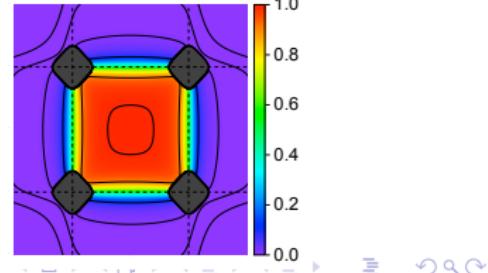
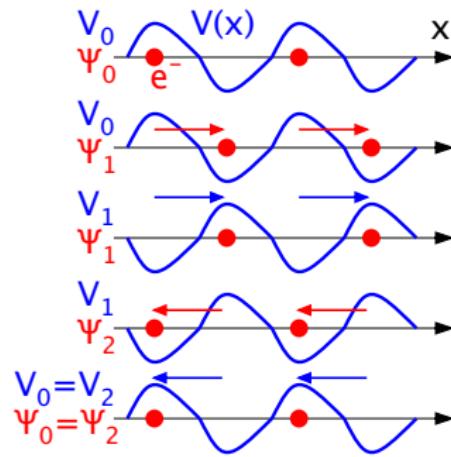
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In spite of its "simplicity" : leads to exotic solutions with anisotropic fermi surfaces.

"Cheap" algorithm : great system sizes are available (up to $10^6 e^-$). We can get rid of finite size effects.



Jellium Model (HEG)

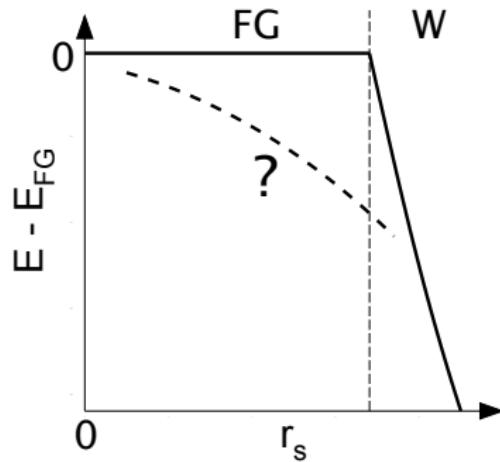
- negative charges : N electrons in a box : $|\Psi_N\rangle$
- positive charges : uniform background
- coulomb interaction : $1/r$

→ good model for solid sodium (Na) (3D), heterostructures (2D)

One physical parameter

Mean distance between electrons $\sim r_s$

- kinetic energy $\sim 1/r_s^2$
- potential energy $\sim 1/r_s$



Two well known solutions :

- Fermi Gas (FG) : minimise kinetic energy
- Wigner crystal (W) : minimise potential energy

Our study : periodic states

Hartree-Fock : $\Psi_N = \det(\{\phi_{\mathbf{k}}\})$

$\phi_{\mathbf{k}}$: single particle wave function.

The goal is to find the set of $\phi_{\mathbf{k}}$ which give the lowest energy.

Ansatz :

Periodic charge density $\rightarrow \phi_{\mathbf{k}} =$ Bloch waves

$$\phi_{\mathbf{k}}(\mathbf{r}) = e^{i\mathbf{k}\mathbf{r}} u_{\mathbf{k}}(\mathbf{r}) \Leftrightarrow |\phi_{\mathbf{k}}\rangle = \sum_{\mathbf{q}} a_{\mathbf{k},\mathbf{q}} |\mathbf{k} + \mathbf{q}\rangle$$

The number of parameters in Ψ_N is greatly reduced. So :

- great system size : up to 10^6 electrons
- good precision on energy (< 0.01 mHa)

Wigner Crystal

Electron density in real space is symmetric.
For example : 2D square lattice.

Bloch waves :

$$\phi_{\mathbf{k}}(\mathbf{r}) = e^{i\mathbf{k}\cdot\mathbf{r}} u_{\mathbf{k}}(\mathbf{r}) \Leftrightarrow |\phi_{\mathbf{k}}\rangle = \sum_{\mathbf{q}} a_{\mathbf{k},\mathbf{q}} |\mathbf{k} + \mathbf{q}\rangle$$

where \mathbf{k} belongs to B , the first Brillouin zone of the lattice.

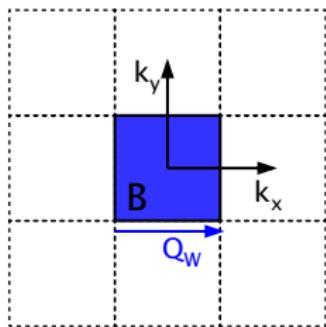
Constraint :

Number of states in B = Number of electrons

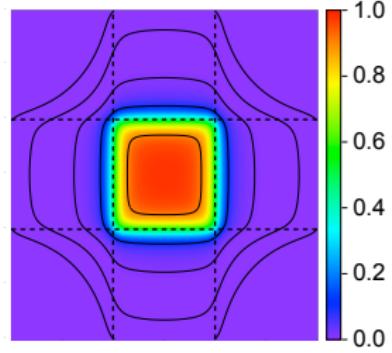
$$Q_W = \alpha_{\text{square}} k_F$$

$$\alpha_{\text{square}} = \sqrt{\pi} \approx 1.77$$

Brillouin zones :



$$n(\mathbf{k})$$



Incommensurate Crystal

Same as Wigner crystal, but :

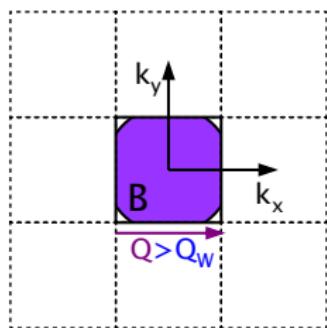
Incommensurate crystal :

Number of states in $B >$ Number of electrons
 $Q > Q_W$

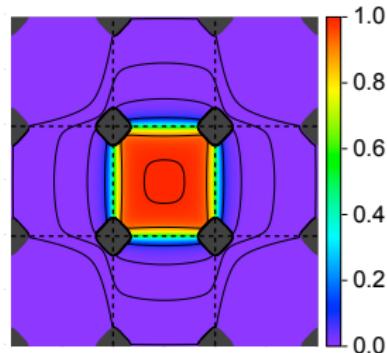
In real space, the electron gas form a crystal with less than one electron per cell.

Note : In the figure on the right, we considered $256^2 \approx 65500$ states in B , so $N \approx 60000$.

Brillouin zones :



$n(\mathbf{k})$



Interpolation between W and FG

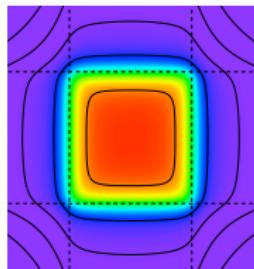
Wigner :

$$Q = Q_W$$

$$Q \approx 1.77k_F$$

$$Q_W$$

$$n(\mathbf{k})$$

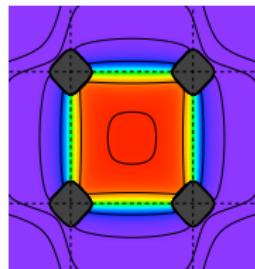


Incommensurate :
 $Q > Q_W$

$$Q \approx 1.84k_F$$

$$Q > Q_W$$

$$n(\mathbf{k})$$

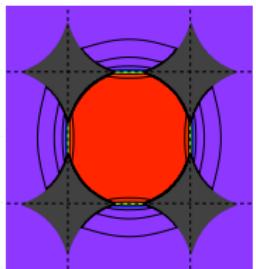


Incommensurate :
 $Q > Q_W$

$$Q \approx 1.98k_F$$

$$Q' > Q$$

$$n(\mathbf{k})$$

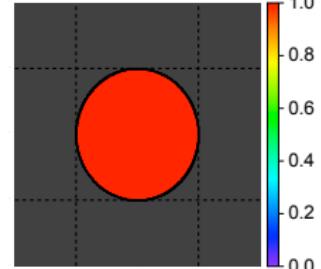


Fermi Gas :

$$Q \geq 2k_F$$

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$$n(\mathbf{k})$$

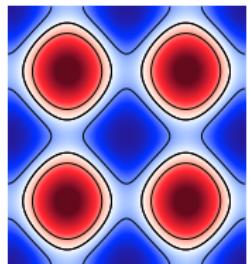


What happens in real space

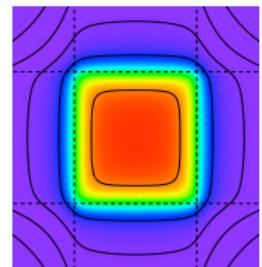
Wigner :

$$Q = Q_W$$
$$Q \approx 1.77k_F$$

$$n(\mathbf{r})$$



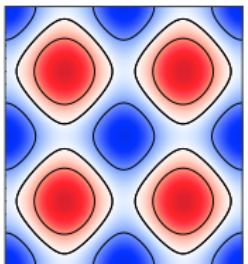
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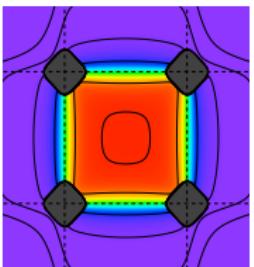
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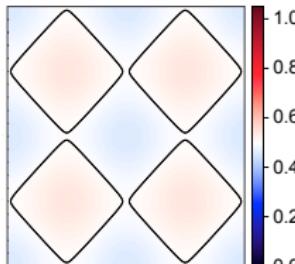
$$n(\mathbf{k})$$



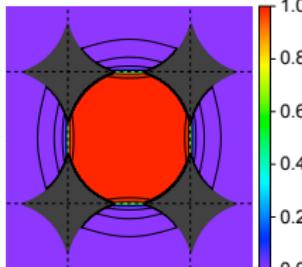
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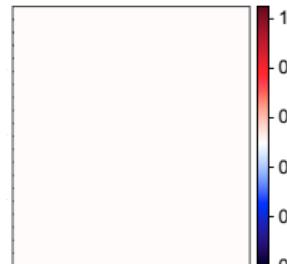
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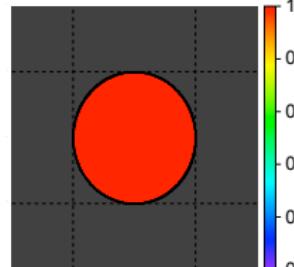
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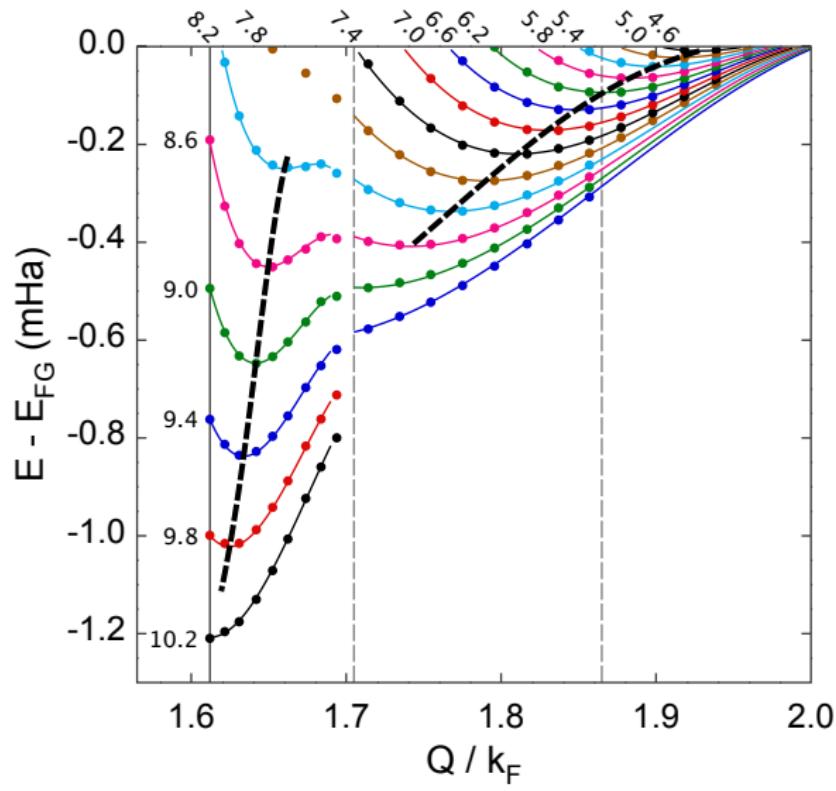


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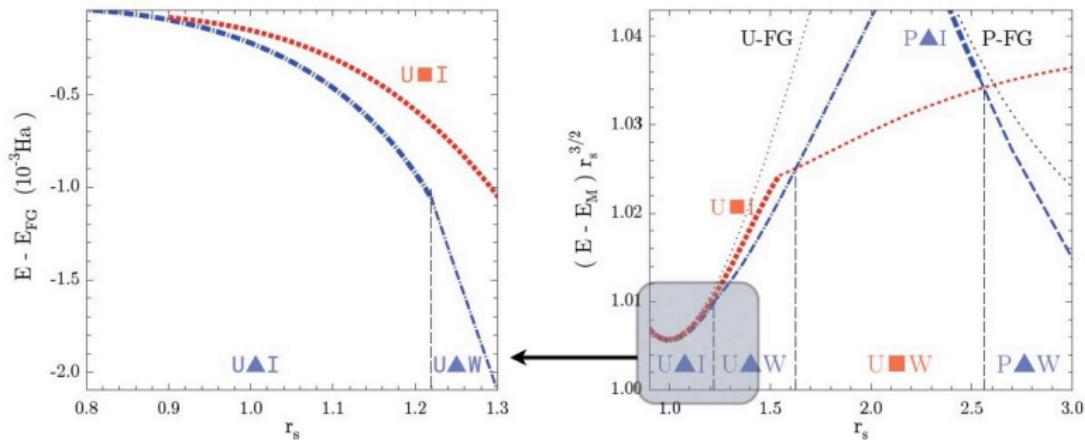
Numerical results at fixed geometry

Example : polarized gas with cubic geometry (3D)



2D phase diagram

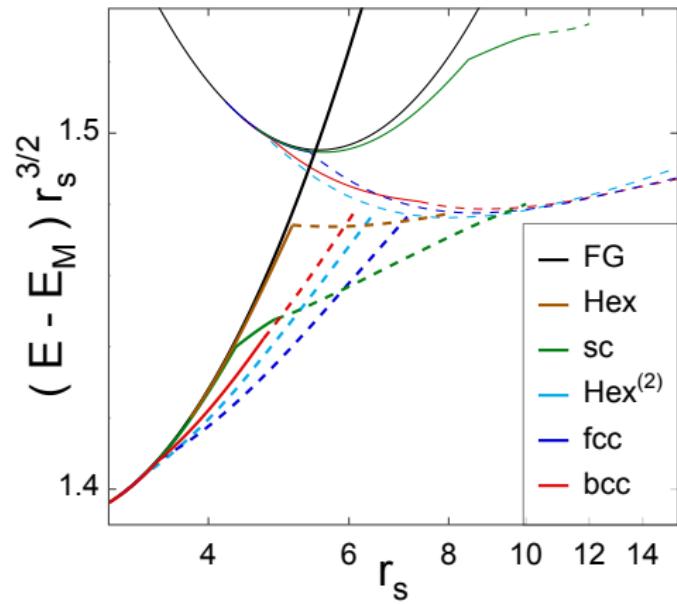
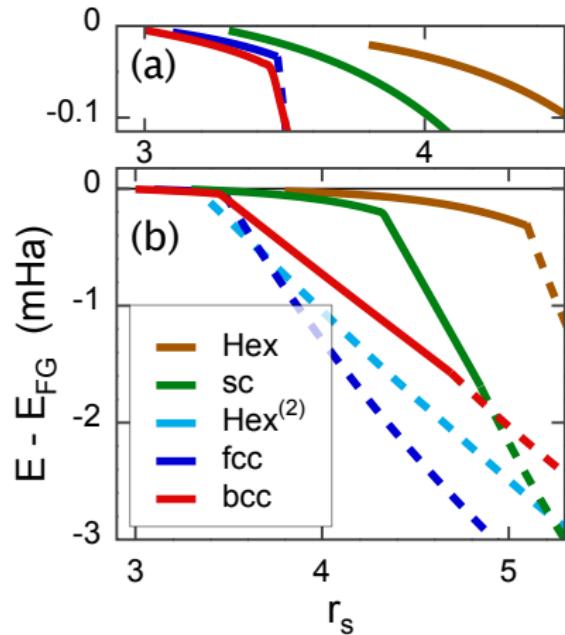
Bernu, Delyon, Holzmann, and Baguet, Phys. Rev. B 84, 115115, 2011



$2.55 < r_s$	\blacktriangle «Wigner»	$\uparrow\uparrow$
$1.55 < r_s < 2.55$	\blacksquare «Wigner»	$\uparrow\downarrow$
$1.225 < r_s < 1.55$	\blacktriangle «Wigner»	$\uparrow\downarrow$
$r_s < 1.225$	\blacktriangle «Incommensurate»	$\uparrow\downarrow$

3D phase diagram

Baguet, Delyon, Bernu, and Holzmann, Phys. Rev. Lett. 111, 166402, 2013



Conclusions

- Transition from the Wigner crystal (large r_s) to the Fermi gas via incommensurate states
- At the thermodynamical limit, the Fermi gas is reached at $r_s = 0$.
- Incommensurate states : crystalline order with a larger modulation than the Wigner crystal ($Q_W < Q(r_s) < 2k_F$).
At small r_s , these states are well described as Spin Density Waves (SDW).

